

# Azimuthal correlations

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# Azimuthal correlations

Why do we want to **measure** (two-particle) **azimuthal correlations** in non-central collisions?

- Resonance **flow**
- Azimuthally sensitive **interferometry**
- **Correlations** between high- $p_T$  particles

Different physical phenomena investigated...

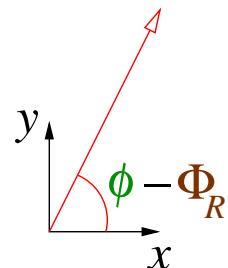
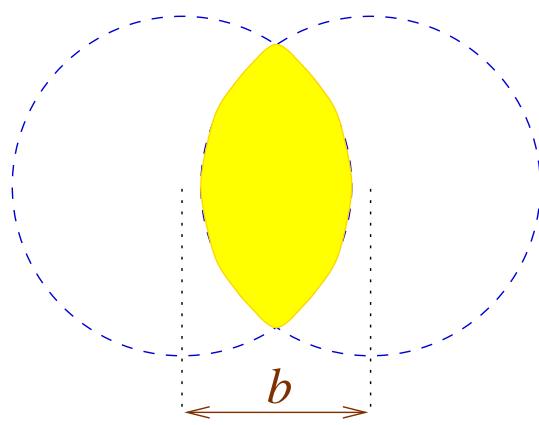
...different methods of analysis

A single **observable**?



# Anisotropic flow

Non-central collision:

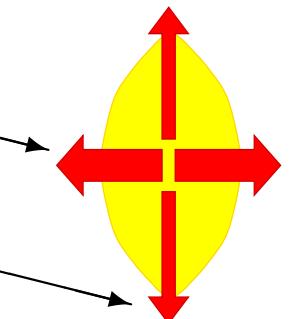


The **particle source** is anisotropic  
(and around it there is only vacuum)

⇒ the **pressure gradient** along the  
**impact parameter** direction  
is stronger than the **gradient**  
perpendicular to the reaction  
plane

⇒ **anisotropic** particle emission: **FLOW**  
in *momentum* space

Particles mainly emitted *in-plane* ( $\phi = \Phi_R$ )  
rather than *out-of-plane* ( $\phi - \Phi_R = 90^\circ$ ).



# Anisotropic flow

Anisotropy quantified by a Fourier expansion:

$$p_1(\phi - \Phi_R) \propto 1 + 2 \textcolor{red}{v}_1 \cos(\phi - \Phi_R) + 2 \textcolor{red}{v}_2 \cos 2(\phi - \Phi_R) + \dots$$

Measuring anisotropic flow is a complicated issue:

Many methods of analysis are available (even unbiased ones!)



# Anisotropic flow

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Measuring anisotropic flow is a complicated issue:

$$v_n = \langle \cos n(\phi - \Phi_R) \rangle$$

lab. frame      not measured!

Many methods of analysis are available (even unbiased ones!)

... at least, for directly detected particles ( $\pi^\pm, K^\pm, p$ )

Improvements still possible for the **flow** of “resonances”

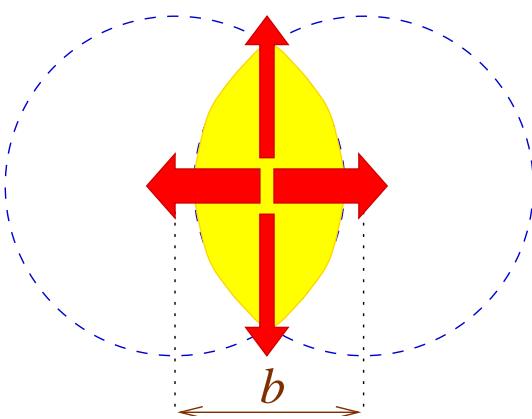
particles that decay before reaching the detector:  $\pi^0$ ,  $K_S^0$ ,  $\Lambda$ ...

⇒ studied through their decay products



# Azimuthally sensitive interferometry

Initial state:

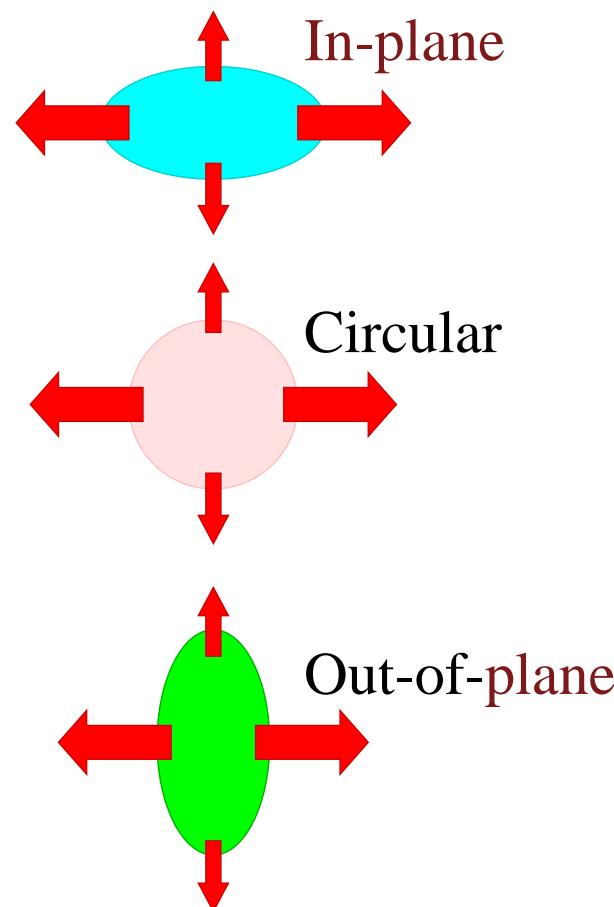


Pressure gradients  
⇒ in-plane flow



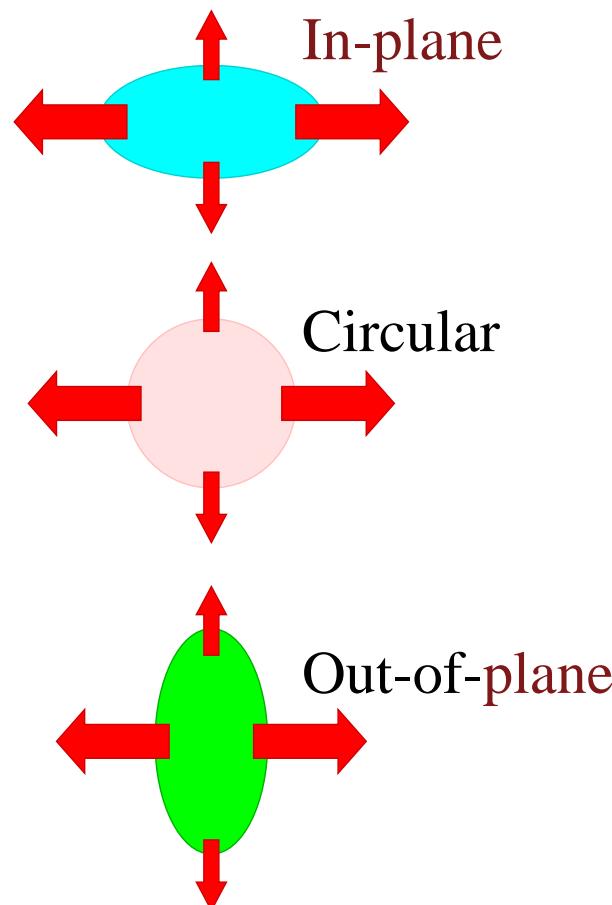
# Azimuthally sensitive interferometry

Final state?



# Azimuthally sensitive interferometry

Final state?



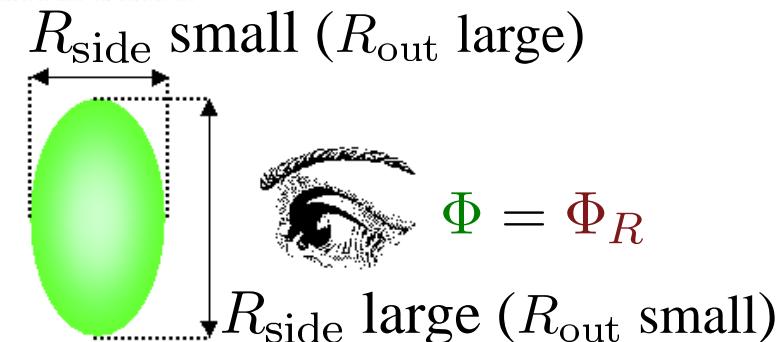
The final space-time configuration can be determined by **two-particle interferometry**

⇒ oscillation of the **HBT radii**  $R_{\text{side}}$ ,  $R_{\text{out}}$   
perpendicular to  $\mathbf{K}_{T\text{pair}}$  →  
along  $\mathbf{K}_{T\text{pair}}$

Example: out-of-plane configuration



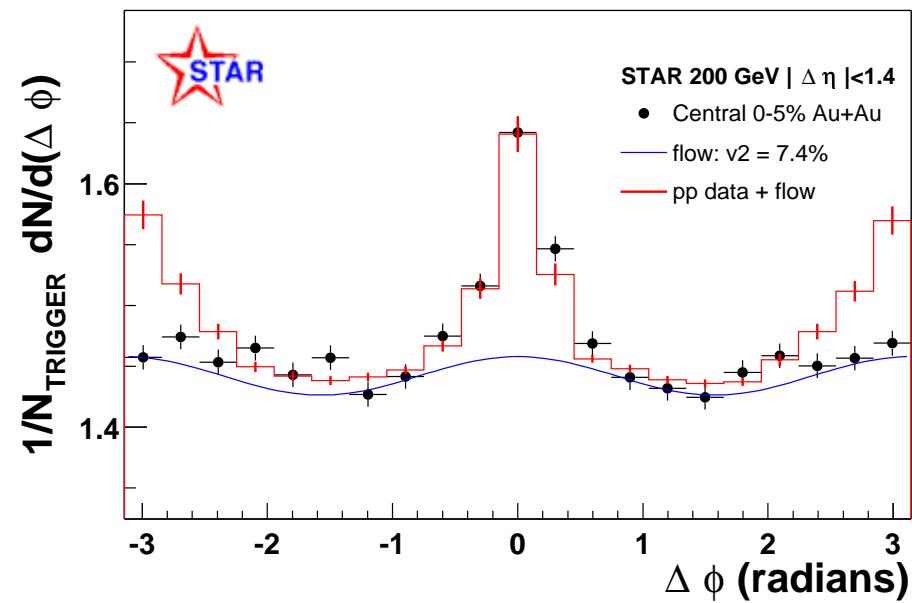
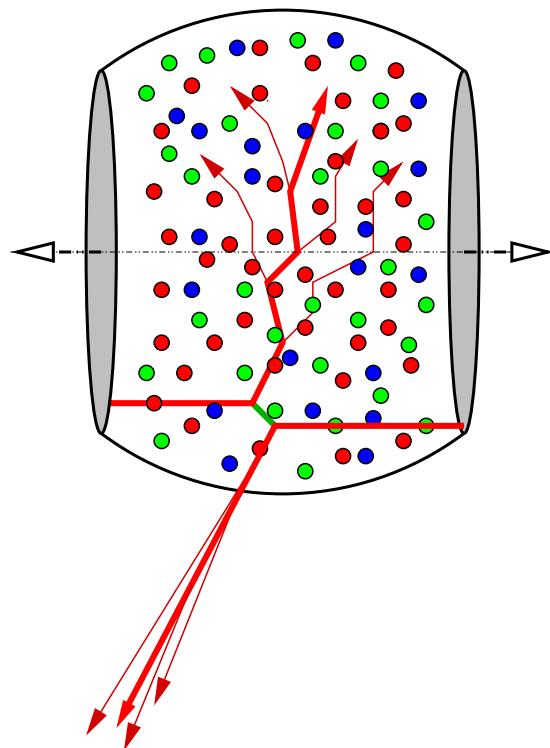
$$\Phi - \Phi_R = 90^\circ$$



# Azimuthal correlations of high $p_T$ particles

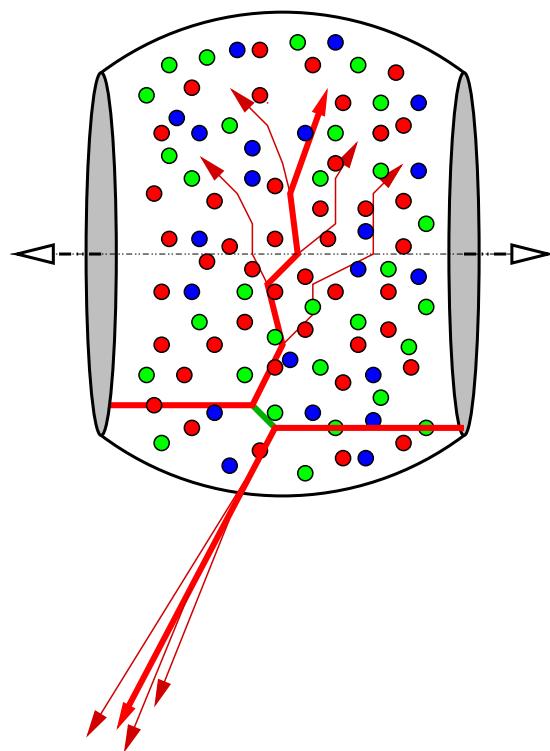
Particles with high momentum (jets) lose energy while traversing the **dense medium** created

⇒ For a pair of jets created close to the edge, only one jet emerges (= detected with high  $p_T$ ), while the back jet is **quenched** (= “not detected”)



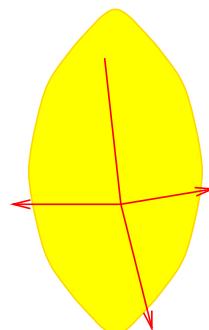
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Amount of **quenching** depends on the length of the **jet path in-medium**



In non-central collisions, **suppression pattern** depends on the **azimuths** of the **high  $p_T$  particles** with respect to the **reaction plane**:

⇒ less **quenching** in-plane ( $\phi = \Phi_R$ ), more out-of-plane



# Two-particle correlations

Two particles characterized by  $(\mathbf{p}_{T1}, \mathbf{p}_{T2}, y_1, y_2), \phi_1, \phi_2$

Rather than  $\begin{cases} \phi_1 - \Phi_R \\ \phi_2 - \Phi_R \end{cases}$ , use  $\begin{cases} \Phi \equiv x\phi_1 + (1-x)\phi_2 : \text{pair azimuth} \\ \Delta\phi \equiv \phi_2 - \phi_1 : \text{relative angle} \end{cases}$

Two-particle distribution:  $p_2(\Phi - \Phi_R, \Delta\phi)$  instead of  $p_2(\phi_1 - \Phi_R, \phi_2 - \Phi_R)$

Fix  $\Delta\phi$ :  $p_2(\Phi - \Phi_R)$  is the pair-angle  $\Phi$  azimuthal distribution, which quantifies the **anisotropy** in  $\Phi - \Phi_R$

... reminiscent of  $p_1(\phi - \Phi_R)$



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Introduce a Fourier expansion!

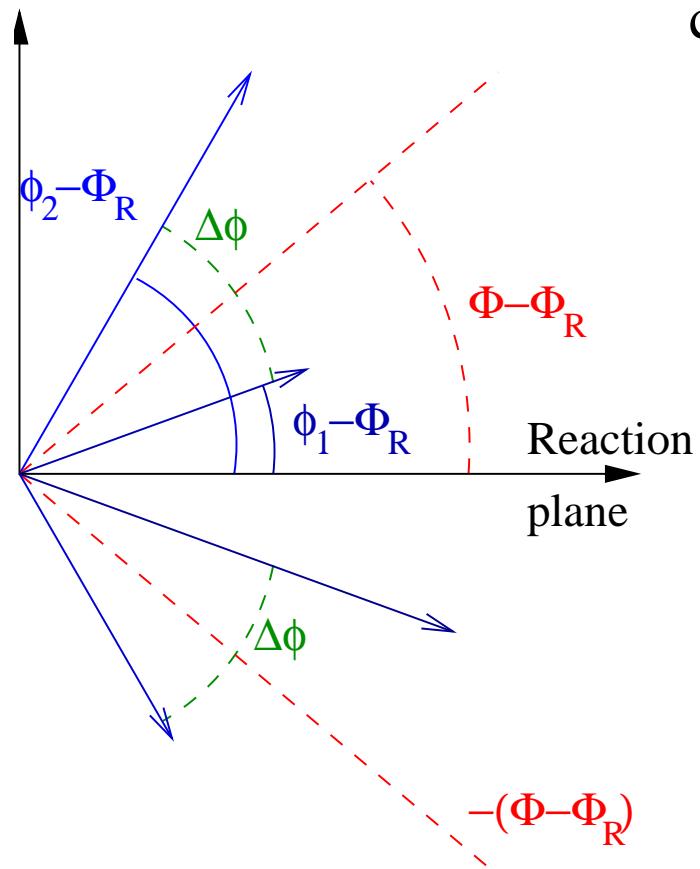
$$p_2(\Phi - \Phi_R) \propto 1 + \sum_{n \neq 0} v_n^{\text{pair}} e^{in(\Phi - \Phi_R)}$$



# “Pair anisotropic flow”

$$p_2(\Phi - \Phi_R) \propto 1 + \sum_{n \neq 0} v_n^{\text{pair}} e^{in(\Phi - \Phi_R)}, \quad v_n^{\text{pair}} = \langle e^{-in(\Phi - \Phi_R)} \rangle$$

cannot be replaced by cos!



Transforming  $\Phi - \Phi_R \rightarrow -(\Phi - \Phi_R)$  with  $\Delta\phi$  constant is *not* a symmetry!

$v_n^{\text{pair}}$  is complex

(≠ “single-particle” anisotropic flow  $v_n$ )



# “Pair anisotropic flow”

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$$= 1 + 2 \sum_{n \geq 1} (v_{c,n}^{\text{pair}} \cos n(\Phi - \Phi_R) + v_{s,n}^{\text{pair}} \sin n(\Phi - \Phi_R))$$

with  $v_{c,n}^{\text{pair}} = \langle \cos(n(\Phi - \Phi_R)) \rangle$ ,  $v_{s,n}^{\text{pair}} = \langle \sin(n(\Phi - \Phi_R)) \rangle$ , real

cf.  $v_n = \langle \cos(n(\phi - \Phi_R)) \rangle$

“new”

Many methods already available for measuring  $v_{c,n}^{\text{pair}}$ !

⇒ all methods of **flow** analysis

... can easily be modified to measure  $v_{s,n}^{\text{pair}}$  (or directly  $v_n^{\text{pair}}$ )



# “Pair anisotropic flow”: jet quenching

$\Phi = \phi_1$  (trigger particle)

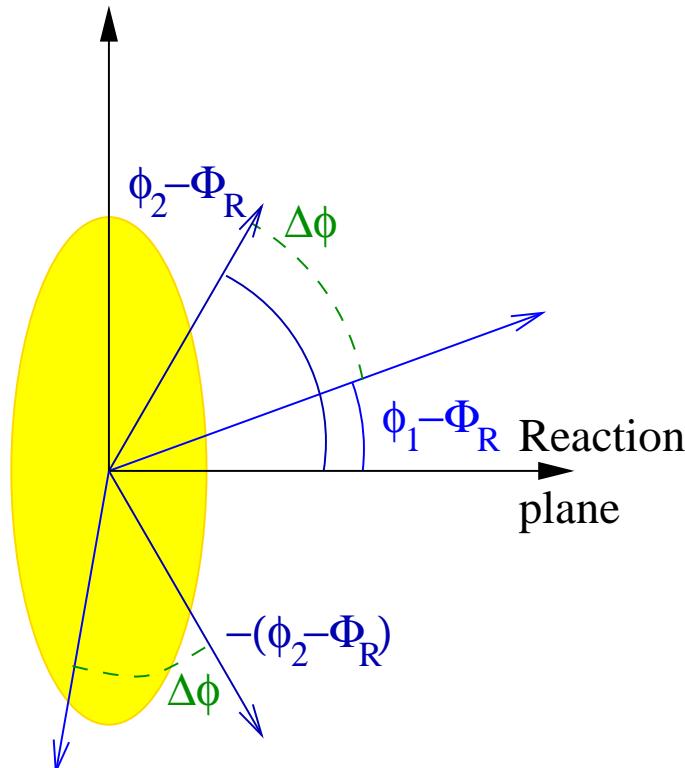
In  $p$ - $p$ , number of associated particles  $\phi_2$  per trigger particle depends on  $\Delta\phi$  only.

If the deficit in high- $p_T$  particles is due to in-medium energy loss, then for a given  $\Delta\phi$  the number of pairs per trigger particle depends on the length of the path followed by the associated particle only

$\Rightarrow$  symmetry  $\phi_2 - \Phi_R \rightarrow -(\phi_2 - \Phi_R)$

⋮

$$v_{s,2}^{\text{pair}}(\Delta\phi) = \left( v_{c,2}^{\text{pair}}(\Delta\phi) - v_2^{\text{trig}} \right) \tan(2\Delta\phi)$$



# Azimuthally sensitive two-particle correlations

Azimuthal dependence of 2-particle correlations = pair **anisotropic flow**

- characterized by pair-**flow** coefficients  $v_{c,n}^{\text{pair}}$ ,  $v_{s,n}^{\text{pair}}$   
model-independent!
- same methods of analysis as for single-particle **anisotropic flow**
  - cumulants, Lee–Yang zeroes:  
no need to estimate  $\Phi_R$ !
- omitted issue: relate  $v_{c,n}^{\text{pair}}$ ,  $v_{s,n}^{\text{pair}}$  to models of the **correlations**:
  - model-dependent predictions
  - different recipes for each specific application (resonance **flow**, **HBT**, high- $p_T$  particle correlations)

